

## Sb<sub>2</sub>O<sub>4</sub> at high pressures and high temperatures

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### Abstract

Investigations on Sb<sub>2</sub>O<sub>4</sub> at high pressure and temperature have been performed up to 600 °C and up to 27.3 GPa. The so-called “high temperature” phase ( $\beta$ -Sb<sub>2</sub>O<sub>4</sub>) was obtained following pressure increase at ambient temperature and at relatively low temperatures. Thus, in contrast to previous perceptions,  $\beta$ -Sb<sub>2</sub>O<sub>4</sub> is the modification more stable at high pressures, i.e., at low temperatures. The fact that the metastable  $\alpha$ -form is typically obtained through the conventional way of preparation has to be attributed to kinetic effects. The pressure-induced phase transitions have been monitored by in-situ X-ray diffraction in a diamond anvil cell, and confirmed ex-situ, by X-ray diffraction at ambient conditions, following temperature decrease and decompression in large volume devices. Bulk modulus values have been derived from the pressure-induced volume changes at room temperature, and are 143 GPa for  $\alpha$ -Sb<sub>2</sub>O<sub>4</sub> and 105 GPa for the  $\beta$ -Sb<sub>2</sub>O<sub>4</sub>.

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### 1. Introduction

The oxides and fluorides of the 14th, 15th and 16th Groups are little known at elevated pressures and temperatures. The common characteristic of these materials is the presence of the non-bonding lone-pair electrons, in many of their binary and tertiary compounds. The reports providing information on the structure, and on the physical and transport properties of such compounds at elevated pressures are rather scarce [1–3], mainly due to their high reactivity. Theoretical models suggest that new electronic properties could be achieved in these structures by applying elevated pressures. In addition, it is long being held that the transformation to a higher density is likely associated with an electronic transition to a metallic behaviour. This is induced by either delocalization of the

lone-pair, or by an overlap of the conduction band with the valence band containing the non-bonding pair.

The present report is part of a larger project dedicated to studying the pressure- ( $p$ ) and temperature ( $T$ )-induced mobility of the electron lone-pairs in oxides and fluorides of elements from the 14th, 15th and 16th Groups [4,5]. We have investigated the Sb–O system at elevated pressure and temperature conditions [6]. Four different binary antimony oxides are known to exist at 1 atm and temperatures up to 1250 °C: Sb<sub>2</sub>O<sub>3</sub>, Sb<sub>2</sub>O<sub>4</sub>, Sb<sub>6</sub>O<sub>13</sub> and Sb<sub>2</sub>O<sub>5</sub> [7]. Literature provides scarce high pressure and high temperature (HP–HT) reports on the Sb–O system, and no data for Sb<sub>2</sub>O<sub>4</sub> were found. Therefore, in the present study we report for the first time high-pressure data on the Sb<sub>2</sub>O<sub>4</sub> system. In practical applications, Sb<sub>2</sub>O<sub>4</sub> is typically used together with other oxides as a catalyst for hydrocarbon cracking and hydrogen reduction, and in the production of ceramic enamels.

The mixed valence Sb<sub>2</sub>O<sub>4</sub> contains Sb in the oxidation states +3 and +5, where Sb(III) has a stereochemically

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active lone electron pair. According to literature data,  $\text{Sb}_2\text{O}_4$  exists in two stable modifications: the room temperature  $\alpha\text{-Sb}_2\text{O}_4$  [8–12] is orthorhombic  $Pna2_1$  (no. 33), with  $a = 5.4426(6)\text{ \AA}$ ,  $b = 4.8098(5)\text{ \AA}$ ,  $c = 11.7830(13)\text{ \AA}$ ,  $V = 308.45(9)\text{ \AA}^3$  and  $Z = 4$ , while the high temperature  $\beta\text{-Sb}_2\text{O}_4$  [10–13] is monoclinic  $C2/c$  (no. 15), with  $a = 12.0582(1)\text{ \AA}$ ,  $b = 4.8344(0)\text{ \AA}$ ,  $c = 5.3827(0)\text{ \AA}$ ,  $\beta = 104.6^\circ$ ,  $V = 303.65(0)\text{ \AA}^3$  and  $Z = 4$ . At 1 atm, the  $\alpha$  to  $\beta$  transition is reported to occur at  $935^\circ\text{C}$  [10,14].

## 2. Experimental details

### 2.1. Large-volume experiments

The starting sample material of crystalline  $\alpha\text{-Sb}_2\text{O}_4$  powder was prepared by heating high-purity  $\text{Sb}_2\text{O}_3$  (Ventron–Alfa Products, ultrapure) in oxygen-flow at  $530^\circ\text{C}$  for 24 h. The obtained powder was ground down, and the whole procedure repeated. The X-ray diffraction patterns obtained (STOE Stadi P instrument with  $\text{CuK}\alpha_1$  radiation, with a curved Johansson-type monochromator using Ge (111), generating a beam with  $1.540598\text{ \AA}$  wavelength), showed only reflections corresponding to  $\alpha\text{-Sb}_2\text{O}_4$ .

For the HP–HT experiments, the antimony oxide was packed in tightly closed 4 mm and 2 mm diameter crucibles. Experiments have been performed using a belt-type press (Dieffenbacher GmbH, Germany), and a multianvil press (Voggenreiter GmbH, Germany) equipped with a Walker module. Both presses have been previously  $p$ - and  $T$ -calibrated. Temperature was measured with Ni/Cr–Ni thermocouples in the belt runs and with W3Re/W26Re thermocouples in the Walker module. After the experiments, X-ray diffraction patterns of the samples were taken at ambient conditions. All samples containing small amounts of sample holder impurities have been removed from the final analysis.

### 2.2. Diamond-anvil cell (DAC) experiments

High-pressure synchrotron X-ray experiments were carried out by using a DAC. The  $T$  301 stainless steel used as gasket was pre-indented to a thickness of  $40\text{ }\mu\text{m}$  from  $250\text{ }\mu\text{m}$ , and then a  $100\text{ }\mu\text{m}$  diameter hole was drilled in its centre by spark erosion. A thin pre-pressed pellet of powder  $\alpha\text{-Sb}_2\text{O}_4$  with thickness about  $10\text{ }\mu\text{m}$  was loaded into the gasket. Silicone oil, which is chemically inert and could generate quasi-hydrostatic pressure conditions, was used as pressure transmitting medium. Two tiny ruby balls were loaded into the sample chamber as pressure markers for the ruby fluorescence method. The pressure values at each step were determined by averaging the pressure readings

measured before and after taking X-ray diffraction patterns. In-situ high-pressure angle-dispersive X-ray diffraction (ADXRD) experiments were performed at room temperature at ID-B station of HPCAT, Advanced Photon Source, Argonne National Laboratory. A double-crystal branching monochromator using water-cooled diamond (111) and Si (220), generated monochromatic beam with wavelength  $\lambda = 0.4157\text{ \AA}$ , focussed by multi-electrode modular piezoelectric bimorph mirrors in Kirkpatrick-Baez configurations down to  $15\text{ }\mu\text{m}$  diameter. A Mar345 image plate with  $100\text{ }\mu\text{m}$  pixel size resolution collected the diffraction patterns, with exposure times between 30 and 150 s. The distance between sample and detector was about 428 mm. Fit2D software [15] was used to integrate the diffraction image. The error in  $\delta d/d$  measurement is estimated to be about  $10^{-4}$ , based on the calibration of silicon standard purchased from the National Institute of Standards and Technology.

## 3. Experimental results

Following the high-pressure runs in large volume presses, the complete transformation of  $\alpha\text{-Sb}_2\text{O}_4$  into  $\beta\text{-Sb}_2\text{O}_4$  was observed at 6 GPa and  $375^\circ\text{C}$ , as seen in Fig. 1. At 4 and 6 GPa, and temperatures between 23 and  $350^\circ\text{C}$ , we noticed the presence of a mixture of  $\alpha\text{-Sb}_2\text{O}_4$  and  $\beta\text{-Sb}_2\text{O}_4$ , with a decrease in the quantitative ratio between the two phases induced by temperature increase. At 6 GPa and  $400^\circ\text{C}$  the diffraction pattern indicated  $\beta\text{-Sb}_2\text{O}_4$  and a small amount of  $\text{Sb}_2\text{O}_3$ .

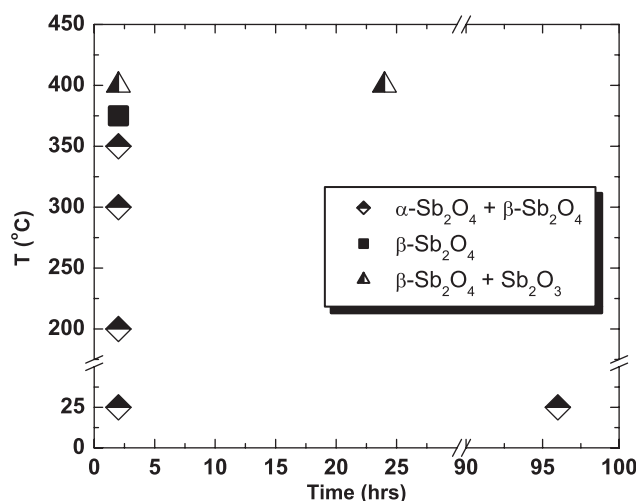


Fig. 1. Result of the high-pressure runs on  $\alpha\text{-Sb}_2\text{O}_4$ , performed at 6 GPa in large volume belt and multianvil presses. The  $\beta\text{-Sb}_2\text{O}_4$  seem to be already occurring at 6 GPa and RT, in the 2 h run, indicating that the fast  $P$ -increase induces the transformation into the more stable phase.

In addition, following runs at 6 GPa and 23 °C, X-ray diffraction analyses revealed the presence of monoclinic  $\beta$ -Sb<sub>2</sub>O<sub>4</sub>, as shown in Fig. 1. The transformation of the  $\alpha$ -phase into  $\beta$ -phase at 6 GPa and ambient temperature is time dependent, but it has a relatively slow kinetics: the amount of the  $\beta$ -phase increased substantially after 24 h, while only traces were found after the 2 h run.

In the DAC experiments performed at 23 °C, the  $\alpha$ -Sb<sub>2</sub>O<sub>4</sub> started transforming at relatively low pressures, but the complete conversion into  $\beta$ -Sb<sub>2</sub>O<sub>4</sub> occurred at pressures much higher than encountered at 375 °C in the large volume presses. The first peaks corresponding to the  $\beta$ -phase appeared at 3.3 GPa (see inset in Fig. 2). Upon further pressure increase, the reflections of the  $\alpha$ -phase decrease gradually. Starting at 20 GPa, only the reflections of the  $\beta$ -phase are visible, and the peaks show a strong pressure-induced broadening. In the X-ray pattern, no more changes are seen until reaching the maximum pressure of 27.3 GPa. On pressure release, no signs of a return to the  $\alpha$ -phase were detected. Furthermore, after 24 h following pressure decrease to

1 atmosphere, only the  $\beta$ -Sb<sub>2</sub>O<sub>4</sub> was detectable in the X-ray pattern.

The DAC experiments provided information regarding the pressure-induced volume changes for the two phases of Sb<sub>2</sub>O<sub>4</sub>. We used the  $p$ – $V$  data to derive the bulk modulus of each individual phase, using the 3<sup>rd</sup>-order formulation of the Birch–Murnaghan Equation of State (EOS) [16]:

$$p_{\text{BM}} = \frac{3}{2} K_{0\text{T}} \left[ \left( \frac{V_0}{V} \right)^{7/3} - \left( \frac{V_0}{V} \right)^{5/3} \right] \times \left\{ 1 - \frac{3}{4} (4 - K'_{0\text{T}}) \left[ \left( \frac{V_0}{V} \right)^{2/3} - 1 \right] \right\}, \quad (1)$$

where  $p$  stands for pressure,  $V_0$  is the volume at  $p = 0$ , and  $K_{0\text{T}}$  and  $K'_{0\text{T}}$  correspond to the bulk modulus and its first pressure derivative, respectively. For this purpose, we employed the EOS fitting program of Angel [17], and we conveniently selected a  $K'_0 = 4$ . The fitted curves are shown in Figs. 3 and 4. The derived

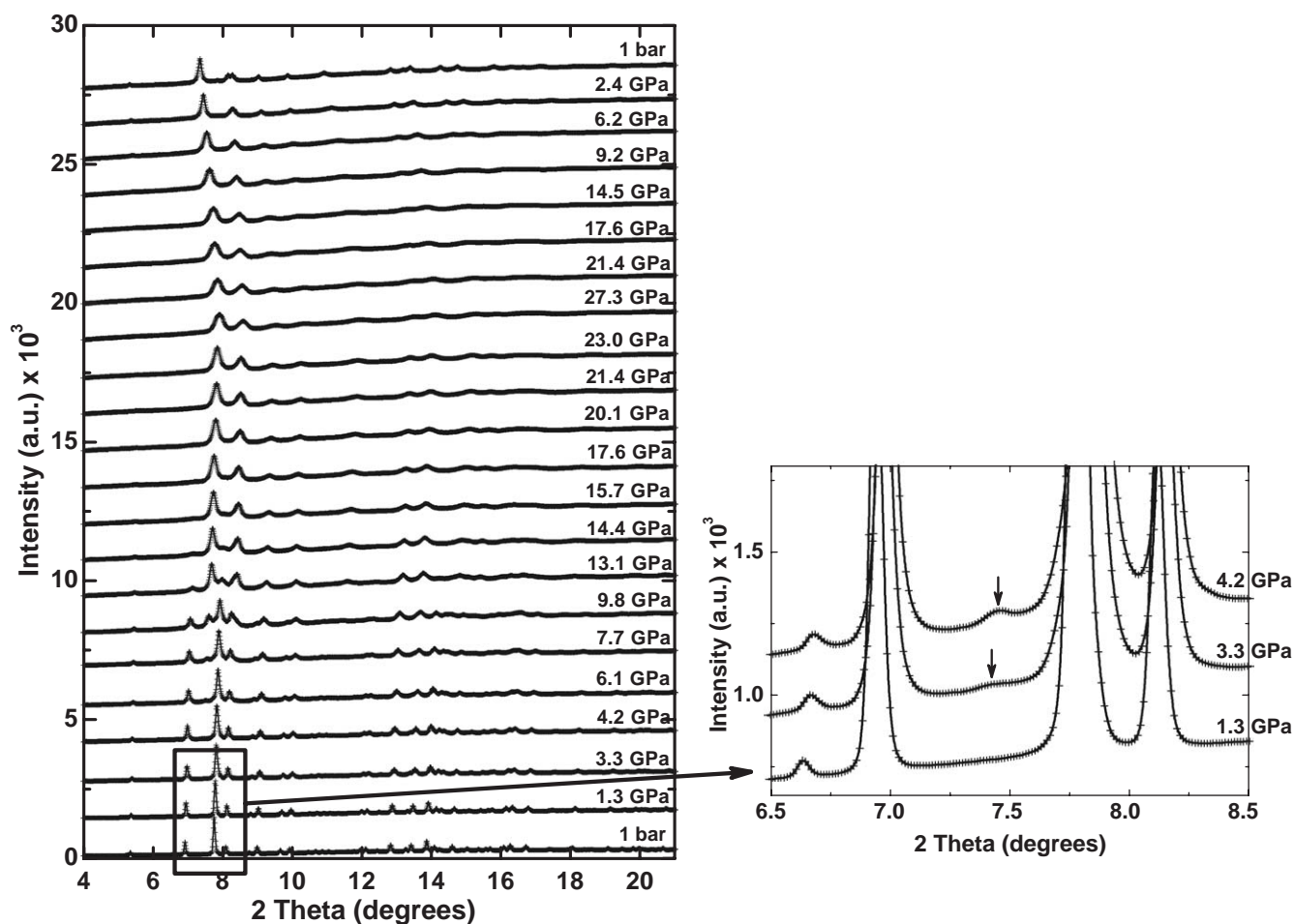


Fig. 2. In-situ X-ray diffraction patterns recorded during pressure increase and decrease at 23 °C in the DAC. The selected region shows the first peaks indicating the occurrence of the  $\beta$ -Sb<sub>2</sub>O<sub>4</sub> phase.

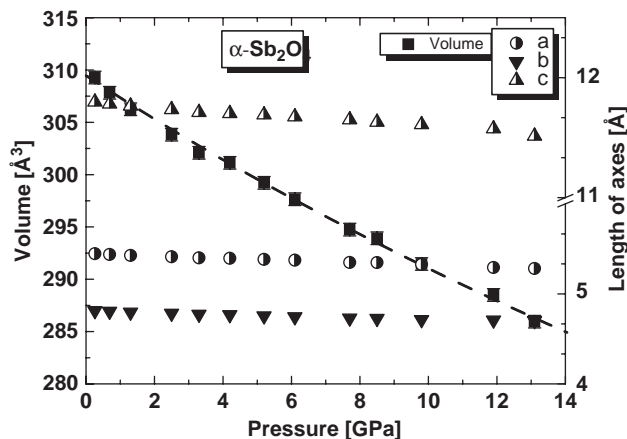


Fig. 3. Pressure-induced changes in lattice size of  $\alpha$ - $\text{Sb}_2\text{O}_4$  during pressure increase (at 23 °C in the DAC), and the derived change in volume. Dashed line is the fitted isothermal Equation of State.

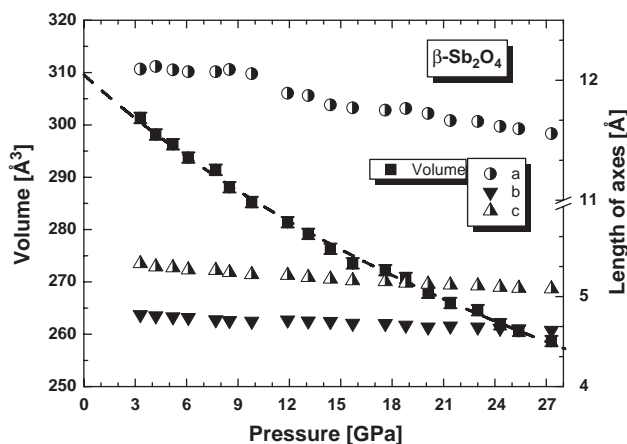


Fig. 4. Pressure-induced changes in lattice size of  $\beta$ - $\text{Sb}_2\text{O}_4$  during pressure increase (at 23 °C in the DAC), and the derived change in volume. Dashed line is the fitted isothermal Equation of State.

bulk modulus values for  $\text{Sb}_2\text{O}_4$  are  $K_{0\alpha} = 143 \pm 1.5$  GPa for the  $\alpha$ -phase, and  $K_{0\beta} = 105 \pm 1.5$  GPa for the  $\beta$ -phase. Similar results were obtained using the Vinet EOS [18]:

$$p_{\text{Vinet}} = \frac{3K_{0T}[1 - (V/V_0)^{1/3}]}{(V/V_0)^{2/3}} \times \exp\left[\frac{3}{2}(K'_{0T} - 1)\left(1 - \left(\frac{V}{V_0}\right)^{1/3}\right)\right]. \quad (2)$$

The relatively low value of  $K_0$  obtained for the  $\beta$ -phase is rather atypical for oxides, most of which usually exhibit an increase in  $K_0$  for high-pressure phases with higher density. However, such anomaly is not unique, being also encountered in elemental Si, Pb and FeS (troilite).

#### 4. Crystal structure of the $\text{Sb}_2\text{O}_4$ phases

The GSAS software [19] was used for crystal structure refinements of both modifications of di-antimony tetraoxide. The starting values for the refinements of the X-ray diffraction patterns of  $\alpha$ - and  $\beta$ - $\text{Sb}_2\text{O}_4$  were taken from the single crystal data by Amador et al. [12]. The refined values for the lattice parameters, the atomic positions and the resulting bond lengths were found to be comparable to previous refinements [9–13]. As shown in Figs. 5 and 6, the crystal structure in both modifications is relatively similar, although  $\alpha$ - $\text{Sb}_2\text{O}_4$  is orthorhombic  $Pna2_1$  and  $\beta$ - $\text{Sb}_2\text{O}_4$  is monoclinic  $C2/c$ . The structure consists, in both phases, of distorted  $\text{Sb}^{\text{V}}\text{O}_6$  octahedra and  $\text{Sb}^{\text{III}}\text{O}_4\text{E}$  pyramids.

The connection between the  $\text{Sb}^{\text{V}}\text{O}_6$  octahedra is established through sharing of corners and through formation of layers, as shown in Figs. 7 and 8. The octahedra around Sb(V) are in the  $\alpha$ -form more distorted than in the  $\beta$ -form. The Sb(III) atoms form

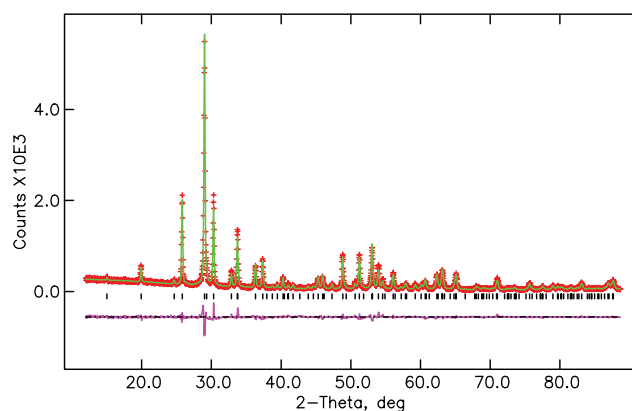


Fig. 5. X-ray diffraction pattern of  $\alpha$ - $\text{Sb}_2\text{O}_4$  at ambient  $p$ ,  $T$  conditions. Crosses stand for measured data; the line is the calculated pattern. The difference is plotted under the bars representing the positions of the Bragg diffraction peaks.

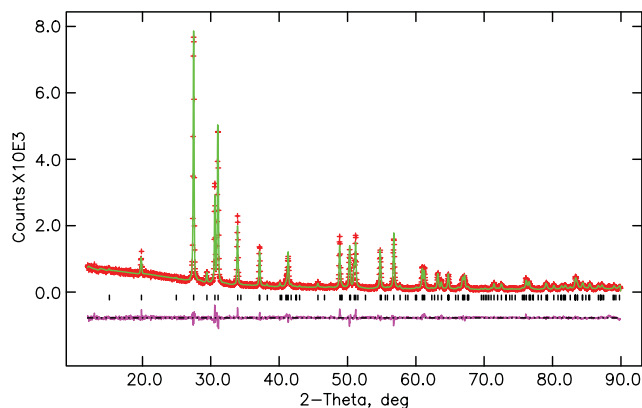
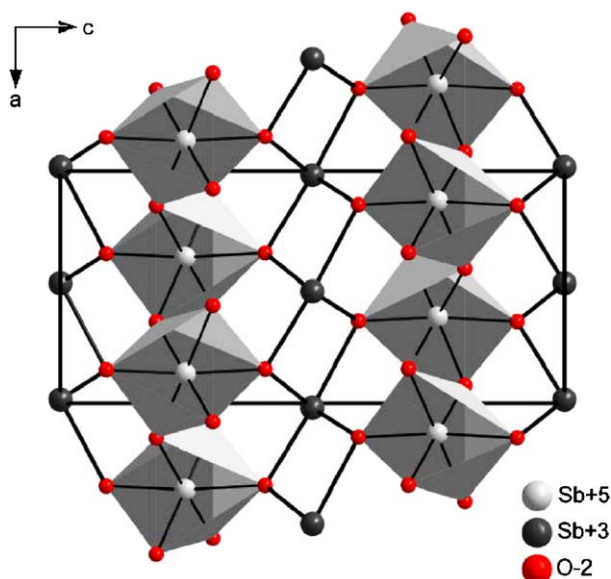
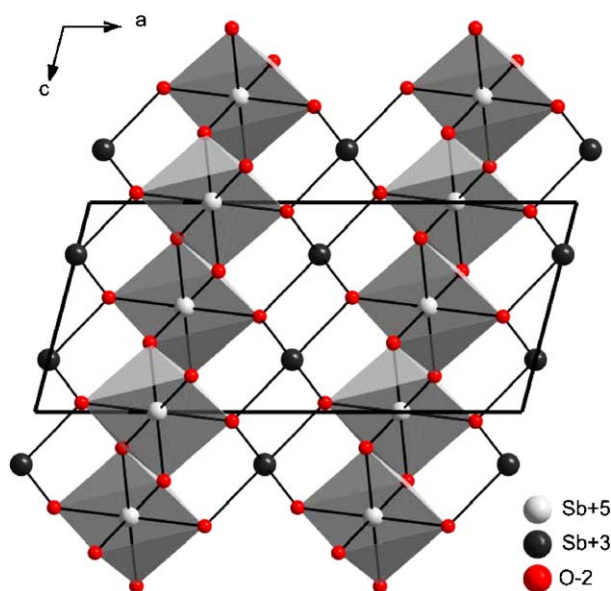


Fig. 6. Ambient conditions X-ray diffraction pattern of  $\beta$ - $\text{Sb}_2\text{O}_4$ , following a run at 6 GPa and 375 °C. Crosses are the experimental data; the line is the calculated pattern. The difference is plotted under the bars representing the positions of the Bragg diffraction peaks.

Fig. 7. Structure of  $\alpha$ - $\text{Sb}_2\text{O}_4$ .Fig. 8. Structure of  $\beta$ - $\text{Sb}_2\text{O}_4$ .

rows parallel to these layers and build up the  $\text{Sb}^{\text{III}}\text{O}_4\text{E}$  units, which share a common edge. Two oxygen atoms belong to the upper layer, and two to the lower layer of the pyramids. The lone electron pair at top of the  $\text{Sb}(\text{III})\text{O}_4\text{E}$  pyramids points always in the opposite direction compared to the neighbouring unit.

## 5. General conclusions

We report here, for the first time, high-pressure and temperature experiments performed on  $\text{Sb}_2\text{O}_4$ . Using both DAC and large volume presses, we determine that

the  $\beta$ -modification of  $\text{Sb}_2\text{O}_4$  is obtained when pressure is applied, under ambient, as well as relatively high temperature. We found the  $\beta$ -phase to be particularly stable at ambient conditions (post quench and decompression), and that the transformation is irreversible, even after heating to temperatures up to  $800^\circ\text{C}$  ( $T$ -dependent XRD-runs). Our crystal structure refinement shows, that at ambient  $p$ - and  $T$ -conditions,  $\beta$ - $\text{Sb}_2\text{O}_4$  has a cell volume of  $V = 303.62 \text{ \AA}^3$  and a density of  $\rho = 6.73 \text{ g/cm}^3$ . By comparison, the values corresponding to  $\alpha$ - $\text{Sb}_2\text{O}_4$ , which is the actual product of the common way of preparation at 1 atm, are  $V = 308.47 \text{ \AA}^3$  and  $\rho = 6.62 \text{ g/cm}^3$ . The result implies that the  $\beta$ -phase is more efficiently packed than  $\alpha$ - $\text{Sb}_2\text{O}_4$ , which until now was assumed to be the more stable one at ambient  $p$  and  $T$ . Typically, a “high-temperature” modification is expected to have a larger volume and a lower density, in comparison to the room temperature phase. However, according to the outcome of our high temperature (large volume presses) and room temperature (DAC) experiments, we conclude that  $\beta$ - $\text{Sb}_2\text{O}_4$  is the thermodynamically stable modification of antimony(III,V) oxide for the entire  $p$ – $T$  range covered by the present study. This finding is also supported by the fact that, under elevated pressures and temperatures,  $\text{Sb}_2\text{O}_3$  is decomposing, forming the  $\beta$ -phase of  $\text{Sb}_2\text{O}_4$  and elemental Sb [6]. Our results validate the suggestions previously made by Keve and Skapski [20], that the  $\beta$ -phase could be the more stable one, and that it is not expected to revert to the  $\alpha$ -phase by heating.

The pressure homology rule [21] specifies that the space group and structure of a high-pressure phase should be similar to the space group of the analogue compound of the next heavier element in the same group at ambient conditions. Bismuth(III,V) oxide, which is monoclinic  $C2/c$  and is a structural analogue of  $\beta$ - $\text{Sb}_2\text{O}_4$ , was synthesized by Kumada et al. [22] at ambient pressure. Apparently, the rule should be applicable for  $\beta$ - $\text{Sb}_2\text{O}_4$  and  $\text{Bi}_2\text{O}_4$ : both have the same structure, and the  $\beta$ -form of  $\text{Sb}_2\text{O}_4$  could be obtained by applying high pressure. However, the  $\alpha$  to  $\beta$  transition is also temperature induced, and the  $\alpha$ - $\text{Sb}_2\text{O}_4$  is not recovered by cooling the  $\beta$ -modification. Since in both cases,  $\text{Sb}_2\text{O}_4$  and  $\text{Bi}_2\text{O}_4$ , the modifications stable at ambient conditions are isostructural, the pressure homology rule does not apply.

## Acknowledgments

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## Appendix A

Tables of the refined crystallographic data (Tables 1 and 2).

Table 1

Room temperature crystallographic and least-square refinement data for  $\alpha$ -Sb<sub>2</sub>O<sub>4</sub> (before *P*-increase, at 0 and 9.8 GPa) and  $\beta$ -Sb<sub>2</sub>O<sub>4</sub> (at 9.2 and 0 GPa, following *P*-decrease)

(a)		
$\alpha$ -Sb <sub>2</sub> O <sub>4</sub>	0 GPa	9.8 GPa
Space group	<i>Pna</i> 2 <sub>1</sub> (no. 33)	<i>Pna</i> 2 <sub>1</sub> (no. 33)
Unit cell dimensions (Å)	<i>a</i> = 5.4524(1) <i>b</i> = 4.8180(1) <i>c</i> = 11.8083(2)	<i>a</i> = 5.3336(17) <i>b</i> = 4.7138(15) <i>c</i> = 11.6146(27)
Cell volume (Å <sup>3</sup> )	<i>V</i> = 310.20(1)	<i>V</i> = 292.01(15)
Z	4	4
Density (g/cm <sup>3</sup> ) <sup>a</sup>	6.584	6.994
Least-square refinement data (%) <sup>b</sup>		
w <i>R</i> <sub>p</sub>	1.17	2.54
<i>R</i> <sub>p</sub>	0.71	1.39
<i>R</i> (F <sup>2</sup> )	5.82	13.12
(b)		
$\beta$ -Sb <sub>2</sub> O <sub>4</sub>	0 GPa	9.2 GPa
Space group	<i>C</i> 2/ <i>c</i> (no. 15)	<i>C</i> 2/ <i>c</i> (no. 15)
Unit cell dimensions (Å)	<i>a</i> = 12.1155(22) <i>b</i> = 4.8465(8) <i>c</i> = 5.4014(11) $\beta$ = 104.59(1)°	<i>a</i> = 11.8960(18) <i>b</i> = 4.7567(15) <i>C</i> = 5.2449(17) $\beta$ = 106.73(3)°
Cell volume (Å <sup>3</sup> )	<i>V</i> = 306.93(10)	<i>V</i> = 284.22(25)
Z	4	4
Density (g/cm <sup>3</sup> ) <sup>a</sup>	6.654	7.186
Least-square refinement data (%) <sup>b</sup>		
w <i>R</i> <sub>p</sub>	0.86	0.56
<i>R</i> <sub>p</sub>	0.51	0.41
<i>R</i> (F <sup>2</sup> )	15.37	11.97

<sup>a</sup>Calculated values

<sup>b</sup>As given in the GSAS manual.

Table 2

Atomic positions and thermal parameters of  $\alpha$ -Sb<sub>2</sub>O<sub>4</sub> and  $\beta$ -Sb<sub>2</sub>O<sub>4</sub>, following Rietveld refinement of DAC data during *P*-increase (for the  $\alpha$ -phase) and *P*-decrease (for the  $\beta$ -phase), for similar pressure ranges. Refined were lattice parameters, atomic positions, parameters of the profile function and background function

Atom	Site <i>X</i>	<i>Y</i>	<i>Z</i>	<i>U</i> <sub>iso</sub> (Å <sup>2</sup> )
$\alpha$ -Sb <sub>2</sub> O <sub>4</sub>	<i>P</i> = 0 GPa/ <i>V</i> = 310.20 Å <sup>3</sup>			
Sb1 (+5)	4 <i>a</i>	0.37408(28)	0.9972(20)	0.2455(6) 0.0005
Sb2 (+3)	4 <i>a</i>	0.9803(8)	0.0352(4)	0.0010(6) 0.0081(9)
O1	4 <i>a</i>	0.338(4)	0.140(4)	0.0879(11) 0.0020
O2	4 <i>a</i>	0.1592(22)	0.6919(23)	0.1883(13) 0.0020

Table 2 (continued)

Atom	Site <i>X</i>	<i>Y</i>	<i>Z</i>	<i>U</i> <sub>iso</sub> (Å <sup>2</sup> )
O3	4 <i>a</i>	0.0708(23)	0.2129(28)	0.3133(11) 0.0020
O4	4 <i>a</i>	0.341(4)	0.806(4)	0.4006(12) 0.0020
$\alpha$ -Sb <sub>2</sub> O <sub>4</sub>	<i>P</i> = 9.8 GPa/ <i>V</i> = 292.01 Å <sup>3</sup>			
Sb1 (+5)	4 <i>a</i>	0.3775(23)	1.006(7)	0.2449(16) 0.043(7)
Sb2 (+3)	4 <i>a</i>	0.9475(31)	0.0530(26)	0.0018(23) 0.045(8)
O1	4 <i>a</i>	0.266(12)	0.147(20)	0.090(4) 0.0020
O2	4 <i>a</i>	0.093(10)	0.791(14)	0.180(4) 0.0020
O3	4 <i>a</i>	0.178(7)	0.318(10)	0.315(4) 0.0020
O4	4 <i>a</i>	0.354(13)	0.833(19)	0.4056(33) 0.0020
$\beta$ -Sb <sub>2</sub> O <sub>4</sub>	<i>P</i> = 0 GPa/ <i>V</i> = 306.93 Å <sup>3</sup>			
Sb1 (+5)	4 <i>c</i> 1/4	1/4	0	0.020(4)
Sb2 (+3)	4 <i>e</i> 0	0.2900(17)	1/4	0.060(7)
O1	8 <i>f</i> 0.1854(18)	−0.052(6)	0.749(6)	0.0020
O2	8 <i>f</i> 0.1067(15)	0.413(5)	1.036(5)	0.0020
$\beta$ -Sb <sub>2</sub> O <sub>4</sub>	<i>P</i> = 9.2 GPa/ <i>V</i> = 284.22 Å <sup>3</sup>			
Sb1 (+5)	4 <i>c</i> 1/4	1/4	0	0.0110(23)
Sb2 (+3)	4 <i>e</i> 0	0.2791(16)	1/4	0.087(5)
O1	8 <i>f</i> 0.1742(12)	−0.052(4)	0.7363(34)	0.0020
O2	8 <i>f</i> 0.0993(11)	0.395(4)	1.0165(34)	0.0020

## References

- [1] U. Häussermann, P. Berastegui, S. Carlson, J. Haines, J.-M. Leger, *Angew. Chem. Int. Ed.* 40 (2001) 4624.
- [2] D. Becker, H. Beck, *Z. Kristallogr.* 219 (2004) 348.
- [3] X. Wang, F. Zhang, I. Loa, K. Syassen, M. Hanfland, Y.-L. Mathis, *Phys. Stat. Sol. B* 241 (2004) 3168.
- [4] D. Orosel, O. Leynaud, P. Balog, M. Jansen, *J. Solid State Chem.* 177 (2004) 1631.
- [5] R. Dinnebier, S. Carlson, M. Hanfland, M. Jansen, *Am. Mineral.* 88 (2003) 996.
- [6] D. Orosel, P. Balog, M. Jansen, New phases of Sb<sub>2</sub>O<sub>3</sub> under high pressures and high temperatures, in preparation.
- [7] T.B. Massalski (Ed.), *Binary Alloy Phase Diagrams*, second ed., vol. 3, ASM International, Materials Park Ohio, 1990, pp. 2912–2913.
- [8] W. Gründer, H. Pätzold, H. Strunz, *Neues Jahrb. Min.* 5 (1962) 93.
- [9] K. Dählström, *Z. Anorg. Allg. Chem.* 239 (1938) 57.
- [10] D. Rogers, A.C. Skapski, *Proc. Chem. Soc.* (1964) 400–401.
- [11] G. Thornton, *Acta Crystallogr. B* 33 (1977) 1271–1273.
- [12] J. Amador, E. Gutierrez Puebla, M.A. Monge, I. Rasines, C. Ruiz Valero, *Inorg. Chem.* 27 (1988) 1367–1370.
- [13] R. Basso, G. Lucchetti, L. Zefiro, A. Palenzona, *Eur. J. Mineral.* 11 (1999) 95–100.
- [14] R.G. Teller, M.R. Antonio, J.F. Brazdil, M. Mehicic, R.K. Grasselli, *Inorg. Chem.* 24 (1985) 3370–3375.
- [15] A.P. Hammersley, S.O. Svensson, M. Hanfland, A.N. Fitch, D. Häussermann, *High Pressure Res.* 14 (1996) 235–248.
- [16] F.D. Murnaghan, *Am. J. Math.* 59 (1937) 235–260.
- [17] R.J. Angel, *EOS-FIT v.5-2*, [www.freeware](http://www.freeware.virginia.edu), Virginia Tech., 2001–2003.
- [18] P. Vinet, J. Ferrante, J.R. Smith, J.H. Rose, *J. Phys. C* 19 (1986) L467–L473.
- [19] C. Larson, R.B. Von Dreele, *General Structure Analysis System (GSAS)*, Los Alamos National Laboratory Report LAUR, 2000, pp. 86–748.
- [20] E.T. Keve, A.C. Skapski, *J. Solid State Chem.* 8 (1973) 159–165.
- [21] A. Neuhaus, *Chimia* 18 (1964) 93–103.
- [22] N. Kumada, N. Kinomura, P.M. Woodward, A.W. Sleight, *J. Solid State Chem.* 116 (1995) 281–285.